

COUNTEREXAMPLES TO L^p COLLAPSING ESTIMATES

XIUMIN DU AND MATEI MACHEDON

ABSTRACT. We show that certain L^2 space-time estimates for generalized density matrices which have been used by several authors in recent years to study equations of BBGKY or Hartree-Fock type, do not have non-trivial $L^p L^q$ generalizations.

1. INTRODUCTION AND MAIN RESULTS

In recent years, effective equations approximating the evolution of a large number of interacting Bosons or Fermions have been studied extensively. The best known example is the celebrated work of Erdős, Schlein and Yau [5], [6].

Since that work, a number of authors have studied the related Gross-Pitaevskii or BBGKY hierarchies, or the Hartree-Fock or Hartree-Fock-Bogoliubov equations, using harmonic analysis techniques and space-time L^2 estimates for a suitable trace density of solutions of the linear Schrödinger equation. We call such estimates “collapsing estimates”, and list several instances, all in 3 space dimensions (thus, $x \in \mathbb{R}^3$, etc.).

If

$$G(t, x, y, z) = e^{\frac{it(\Delta_x + \Delta_y - \Delta_z)}{2}} G_0, \quad (1)$$

then

$$\|\nabla_x G(t, x, x, x)\|_{L^2(dt dx)} \lesssim \|\nabla_x \nabla_y \nabla_z G_0(x, y, z)\|_{L^2(dx dy dz)}. \quad (2)$$

This estimate was used in the study of the Gross-Pitaevskii or BBGKY hierarchies. See [11] (where the estimate originates), as well as [1], [3], [4].

Another related example is: if

$$\Lambda(t, x, y) = e^{\frac{it(\Delta_x + \Delta_y)}{2}} \Lambda_0, \quad (3)$$

Date: February 21, 2020.

1991 Mathematics Subject Classification. 35Q55.

Key words and phrases. Collapsing estimates.

then

$$\|\|\nabla|_x^{1/2}\Lambda(t, x, x)\|\|_{L^2(dt dx)} \lesssim \|\|\nabla|_x^{1/2}|\nabla|_y^{1/2}\Lambda_0(x, y)\|\|_{L^2(dx dy)}. \quad (4)$$

This estimate is useful for the Hartree-Fock-Bogoliubov equations, see [9], [10].

Finally, if

$$\Gamma(t, x, y) = e^{\frac{it(\Delta_x - \Delta_y)}{2}} \Gamma_0, \quad (5)$$

then

$$\|\|\nabla_x|^{\frac{1}{2}}\langle\nabla_x\rangle^{2\epsilon}\Gamma(t, x, x)\|\|_{L^2(dt dx)} \lesssim_\epsilon \|\|\langle\nabla_x\rangle^{\frac{1}{2}+\epsilon}\langle\nabla_y\rangle^{\frac{1}{2}+\epsilon}\Gamma_0(x, y)\|\|_{L^2(dx dy)}. \quad (6)$$

Such estimates are relevant to both the Hartree-Fock-Bogoliubov equations mentioned above, and Hartree-Fock. See Theorem 3.3 in [2].

We also mention the approach of [7], [8] which applies to equation (5) and allows a wide range of $L^p(dt)L^q(dx)$ estimates on the left hand side, but the right hand side of the inequality is estimated in a Schatten norm.

It is natural to ask whether one can replace the $L^2(dt)L^2(dx)$ norm on the left hand side of estimates (2), (4) or (6) by an $L^p(dt)L^q(dx)$ norm, while keeping the right hand side in a Sobolev norm, which is useful for applications to PDEs. One can trivially make p or q bigger than 2 by putting more derivatives on the right hand side, so the interesting question is if one can make p or q less than 2.

The main result of this note is that this is impossible.

We prove the following closely related results.

Theorem 1.1. *Let Λ be given by (3), with $x, y \in \mathbb{R}^n$. Assume*

$$\|\|\nabla|_x^\alpha\Lambda(t, x, x)\|\|_{L^p(dt)L^q(dx)} \lesssim \|\|\Lambda_0(x, y)\|\|_{H^s(dx dy)} \quad (7)$$

for some $\alpha \geq 0, s \geq 0$. Then $p \geq 2$ and $q \geq 2$.

Theorem 1.2. *Let Γ be given by (5), with $x, y \in \mathbb{R}^n$. Assume*

$$\|\|\nabla|_x^\alpha\Gamma(t, x, x)\|\|_{L^p(dt)L^q(dx)} \lesssim \|\|\Gamma_0(x, y)\|\|_{H^s(dx dy)} \quad (8)$$

for some $\alpha \geq 0, s \geq 0$. Then $p \geq 2$ and $q \geq 2$.

Theorem 1.3. *Let G be given by (1), with $x, y, z \in \mathbb{R}^n$. Assume*

$$\|\|\nabla|_x^\alpha G(t, x, x, x)\|\|_{L^p(dt)L^q(dx)} \lesssim \|\|G_0(x, y, z)\|\|_{H^s(dx dy dz)} \quad (9)$$

for some $\alpha \geq 0, s \geq 0$. Then $p \geq 2$ and $q \geq 2$.

Acknowledgements. The first author is supported by the National Science Foundation under Grant No. DMS-1856475.

2. PROOFS

2.1. Proof of Theorem 1.1.

2.1.1. *Necessity of $p \geq 2$.* Let R be a large number (which will approach ∞ at the end of the proof). Let C be a fixed large number (depending on n). Let

$$F_0(x, y) = e^{-\frac{|x|^2 + |y|^2}{2CR}}$$

so that

$$e^{\frac{it(\Delta_x + \Delta_y)}{2}} F_0 := F(t, x, y) = \frac{1}{(1 + it/(CR))^n} e^{-\frac{|x|^2 + |y|^2}{2(CR + it)}}. \quad (10)$$

We think of $F(t, x, y)$ as the basic “vertical tube” solution to the linear Schrödinger equation in $2n + 1$ dimensions which is essentially 1 if $|x|, |y| \leq R^{1/2}$, $0 \leq t \leq R$. The rigorous statement is that C is chosen so that $\Re F(t, x, y) \geq \frac{1}{2}$ in the above range. Also, the Fourier transform (in space) of F is essentially supported at frequencies $|\xi|, |\eta| \leq R^{-1/2}$.

We choose the function $\Lambda(t, x, y)$ to be a sum of translates and modulations of $F(t, x, y)$ which are inclined at 45 degrees and are trained to reach the region $|x| \leq \frac{1}{100}$, $|y| \leq \frac{1}{100}$, $R - R^{\frac{1}{2}} < t < R$ with almost the same oscillation (and almost no cancellations). The summands will have Fourier transforms essentially supported in balls of radius $R^{-1/2}$ centered at unit vectors.

Explicitly, choose roughly $R^{n-\frac{1}{2}}$ points (x_k, y_k) which are spaced at distance $R^{1/2}$ from each other on the sphere $|(x, y)| = R$. For technical reasons, we only choose points for which all coordinates are $\geq \frac{R}{10n}$. Define $(\xi_k, \eta_k) = \frac{(x_k, y_k)}{R}$.

Choose the following initial conditions:

$$\Lambda_0(x, y) = \sum e^{i(x \cdot \xi_k + y \cdot \eta_k)} F_0(x + x_k, y + y_k).$$

The functions being summed are approximately orthogonal and each have L^2 norm $\sim R^{n/2}$:

$$\int |F_0(x + x_k, y + y_k) F_0(x + x_l, y + y_l)| dx dy = \pi^n (CR)^n e^{-\frac{|(x_k, y_k) - (x_l, y_l)|^2}{4CR}}. \quad (11)$$

Recalling that the sum has $\sim R^{n-\frac{1}{2}}$ terms, we derive

$$\|\Lambda_0\|_{L^2(dx dy)} \lesssim R^{n-\frac{1}{4}}.$$

The same type of upper bound holds for higher order derivatives (since $|(\xi_k, \eta_k)| = 1$), thus, for each fixed s ,

$$\|\Lambda_0\|_{H^s(dx dy)} \lesssim R^{n-\frac{1}{4}}. \quad (12)$$

The solution looks like

$$\begin{aligned}\Lambda(t, x, y) &= \sum e^{-it\frac{(|\xi_k|^2 + |\eta_k|^2)}{2}} e^{i(x \cdot \xi_k + y \cdot \eta_k)} F(t, x + x_k - t\xi_k, y + y_k - t\eta_k) \\ &= e^{-i\frac{t}{2}} \sum e^{i(x \cdot \xi_k + y \cdot \eta_k)} F(t, x + x_k - t\xi_k, y + y_k - t\eta_k),\end{aligned}$$

and

$$|\Lambda(t, x, y)| \geq \Re \sum e^{i(x \cdot \xi_k + y \cdot \eta_k)} F(t, x + x_k - t\xi_k, y + y_k - t\eta_k) \sim R^{n-\frac{1}{2}},$$

if $|(x, y)| \leq \frac{1}{100}$, $R - R^{\frac{1}{2}} < t < R$. Thus

$$R^{\frac{1}{2p}} R^{n-\frac{1}{2}} \lesssim \|\Lambda(t, x, x)\|_{L^p(dt)L^q(dx)}, \quad (13)$$

so, recalling (12), if

$$\|\Lambda(t, x, x)\|_{L^p(dt)L^q(dx)} \lesssim \|\Lambda_0(x, y)\|_{H^s(dx dy)},$$

then $p \geq 2$.

Using the product rule and the lower bounds on the components of ξ_k, η_k , same argument works for ordinary derivatives of order $\alpha = m \in \mathbb{N}$.

To justify the statement for fractional derivatives of non-integer order α , do a Littlewood-Paley decomposition in space $\Lambda(t, \cdot, \cdot) = P_{\leq 10}\Lambda(t, \cdot, \cdot) + P_{\geq 10}\Lambda(t, \cdot, \cdot)$, where $P_{\leq 10}$ localizes functions of $2n$ variables, smoothly at frequencies ≤ 10 . Then $P_{\geq 10}\Lambda(t, \cdot, \cdot)$ is exponentially small as $R \rightarrow \infty$. This is true for the function F_0 , and its translates by a unit vector in Fourier space.

A crude estimate is

$$\|P_{\geq 10}\Lambda(t, \cdot, \cdot)\|_{H^s} \lesssim_s e^{-\sqrt{R}}.$$

For our counterexample, we use $P_{\leq 10}\Lambda(t, \cdot, \cdot)$ instead of $\Lambda(t, \cdot, \cdot)$.

Thus, for R sufficiently large, $|\nabla^m P_{\leq 10}\Lambda(t, x, y)| \sim |\nabla^m \Lambda(t, x, y)| \sim R^{n-\frac{1}{2}}$ if $|(x, y)| \leq \frac{1}{100}$, $R - R^{\frac{1}{2}} < t < R$. The function $(P_{\leq 10}\Lambda)(t, x, x)$ is supported, in Fourier space, at frequencies $|\xi| \leq 20$. Denote, by abuse of notation, $P_{\leq 20}$ the operator localizing functions of n variables at frequencies $|\xi| \leq 20$. Let $m \in \mathbb{N}$, $m > \alpha$. Then the operator $\frac{\nabla^m}{|\nabla|^\alpha} P_{\leq 20}$ (defined in the obvious way on the Fourier transform side) is bounded on all L^p spaces, and

$$\begin{aligned}R^{\frac{1}{2p}} R^{n-\frac{1}{2}} &\lesssim \|\nabla^m (P_{\leq 10}\Lambda)(t, x, x)\|_{L^p(dt)L^q(dx)} \\ &= \left\| \frac{\nabla^m}{|\nabla|^\alpha} P_{\leq 20} |\nabla|^\alpha (P_{\leq 10}\Lambda)(t, x, x) \right\|_{L^p(dt)L^q(dx)} \\ &\lesssim \| |\nabla|^\alpha (P_{\leq 10}\Lambda)(t, x, x) \|_{L^p(dt)L^q(dx)},\end{aligned}$$

while

$$\|P_{\leq 10}\Lambda_0\|_{H^s(dx dy)} \lesssim C^n R^{n-\frac{1}{4}}.$$

Letting $R \rightarrow \infty$, we conclude $p \geq 2$ as before.

2.1.2. *Necessity of $q \geq 2$.* Let $F(t, x, y)$ be the basic vertical tube solution of height R (as (10)). Let $m \gg 1$. Choose roughly $R^{mn-\frac{n}{2}}$ points x_k which are spaced at distance $\sim R^{\frac{1}{2}}$ in a large ball $B(0, R^m)$ of radius R^m in \mathbb{R}^n . Fix a unit vector $\xi \in S^{n-1}$.

We take initial conditions

$$\Lambda_0(x, y) = e^{i(x+y)\cdot\xi} \sum F_0(x + x_k, y + x_k).$$

Then

$$\Lambda(t, x, y) = e^{i(x+y)\cdot\xi} e^{-it} \sum F(t, x + x_k - t\xi, y + x_k - t\xi).$$

There are roughly $R^{mn-\frac{n}{2}}$ terms in the sum. The summands are essentially orthogonal (as in (11)) and each term has L^2 norm $\sim R^{n/2}$, thus

$$\|\Lambda_0\|_{L^2(dx dy)} \sim R^{\frac{n}{4} + \frac{mn}{2}}.$$

On the other hand, each $F(t, x + x_k - t\xi, y + x_k - t\xi)$ is essentially 1 on a tube T_k of radius $R^{1/2}$ and length R in $2n+1$ dimensions, and rapidly decaying out of T_k . Note that at $t = 0$, T_k is centered at $(0, -x_k, -x_k)$. Moreover, these tubes T_k are in the same direction $(1, \xi, \xi)$ and hence disjoint. Therefore, $|\Lambda(t, x, y)| \gtrsim 1$ on the union of the tubes T_k . In particular, $|\Lambda(t, x, x)| \gtrsim 1$ for $0 \leq t \leq R$ and $x \in B(t\xi, R^m)$. We only need the previous estimate for $0 \leq t \leq 1$, where the claim is obvious. In addition, the Fourier transform of $\Lambda(t, x, x)$ is supported (essentially) in a $R^{-\frac{1}{2}}$ neighbourhood of the point 2ξ , with $|\xi| = 1$, so $\|\nabla|^\alpha \Lambda(t, x, x)\| \gtrsim 1$ for $0 \leq t \leq 1$ and $x \in B(t\xi, R^m)$. Thus

$$\|\nabla|^\alpha \Lambda(t, x, x)\|_{L^p([0,1])L^q(dx)} \gtrsim R^{\frac{mn}{q}},$$

while $\|\Lambda_0\|_{H^s(dx dy)} \sim \|\Lambda_0\|_{L^2(dx dy)} \sim R^{\frac{n}{4} + \frac{mn}{2}}$ and $m \gg 1$, so $q \geq 2$ is necessary.

2.2. Proof of Theorem 1.2. The examples for Γ are similar to those for Λ , and are included for completeness.

2.2.1. *Necessity of $p \geq 2$.* First we take the basic “vertical tube” solution. Let

$$F_0(x, y) = e^{-\frac{|x|^2 + |y|^2}{2CR}}$$

so that

$$e^{\frac{it(\Delta_x - \Delta_y)}{2}} F_0 := F(t, x, y) = \frac{1}{(1 + (\frac{t}{CR})^2)^{\frac{n}{2}}} e^{-\frac{|x|^2}{2(CR+it)}} e^{-\frac{|y|^2}{2(CR-it)}}. \quad (14)$$

The solution $F(t, x, y)$ is essentially 1 if $|x|, |y| \leq R^{1/2}$, $0 \leq t \leq R$. More precisely, we choose a large constant $C = C(n)$ so that $\Re F(t, x, y) \geq \frac{1}{2}$ in the above range. Also, as before, the Fourier transform (in space) of F is essentially supported at frequencies $|\xi|, |\eta| \leq R^{-1/2}$.

Pick roughly $R^{n-\frac{1}{2}}$ points (x_k, y_k) which are spaced at distance $\sim R^{1/2}$ from each other on the surface $\{(x, y) : |x| = |y|, \frac{R}{2} \leq |x| \leq R\}$. Define $(\xi_k, \eta_k) = \frac{1}{R}(x_k, y_k)$ so that $|\xi_k|^2 - |\eta_k|^2 = 0$ and $|(\xi_k, \eta_k)| \sim 1$.

Take the following initial conditions

$$\Gamma_0(x, y) = \sum e^{i(x \cdot \xi_k - y \cdot \eta_k)} F_0(x + x_k, y + y_k)$$

so that the solution is

$$\begin{aligned} \Gamma(t, x, y) &= \sum e^{-it \frac{(|\xi_k|^2 - |\eta_k|^2)}{2}} e^{i(x \cdot \xi_k - y \cdot \eta_k)} F(t, x + x_k - t\xi_k, y + y_k - t\eta_k) \\ &= \sum e^{i(x \cdot \xi_k - y \cdot \eta_k)} F(t, x + x_k - t\xi_k, y + y_k - t\eta_k). \end{aligned}$$

Since the $\sim R^{n-\frac{1}{2}}$ terms in Γ_0 are essentially orthogonal and each have L^2 norm $\sim R^{n/2}$, we get

$$\|\Gamma_0\|_{L^2(dx dy)} \lesssim R^{n-\frac{1}{4}}.$$

Moreover, since $|(\xi_k, \eta_k)| \sim 1$, there also holds

$$\|\Gamma_0\|_{H^s(dx dy)} \lesssim R^{n-\frac{1}{4}}. \quad (15)$$

From the expression of Γ , we see that

$$|\Gamma(t, x, y)| \gtrsim R^{n-\frac{1}{2}} \quad \text{for } |(x, y)| \leq \frac{1}{100}, R - R^{\frac{1}{2}} < t < R.$$

Therefore,

$$\|\Gamma(t, x, x)\|_{L^p(dt) L^q(dx)} \gtrsim R^{\frac{1}{2p}} R^{n-\frac{1}{2}},$$

so, recalling (15), if

$$\|\Gamma(t, x, x)\|_{L^p(dt) L^q(dx)} \lesssim \|\Gamma_0(x, y)\|_{H^s(dx dy)},$$

then $p \geq 2$. From a similar argument to the one in subsection 2.1.1 (i.e. only using x_k, y_k for which all coordinates of ξ_k and $-\eta_k$ are $\geq \frac{1}{10n}$), $p \geq 2$ is also necessary for estimates of the form

$$\|\nabla_x^\alpha \Gamma(t, x, x)\|_{L^p(dt)L^q(dx)} \lesssim \|\Gamma_0(x, y)\|_{H^s(dx dy)}.$$

2.2.2. Necessity of $q \geq 2$. Let $F(t, x, y)$ be the basic vertical tube solution of height R (as (14)). Let $m \gg 1$. Choose roughly $R^{mn - \frac{n}{2}}$ points x_k which are spaced at distance $\sim R^{\frac{1}{2}}$ in a large ball $B(0, R^m)$ of radius R^m in \mathbb{R}^n . Fix a unit vector $\xi \in S^{n-1}$.

We take initial conditions

$$\Gamma_0(x, y) = e^{ix \cdot \xi} \sum F_0(x + x_k, y + x_k),$$

so that the solution is

$$\Gamma(t, x, y) = e^{ix \cdot \xi} \sum F(t, x + x_k - t\xi, y + x_k).$$

Note that $\Gamma(t, x, x) \gtrsim 1$ for $0 \leq t \leq 1$ and $|x| \leq R^m$. Moreover, the Fourier transform of $\Gamma(t, x, x)$ is essentially supported in a $R^{-1/2}$ neighborhood of the point ξ with $|\xi| = 1$.

Then, the necessity of $q \geq 2$ follows from the same calculation as in subsection 2.1.2.

2.3. Proof of Theorem 1.3. The examples for G are similar to those in previous subsections.

2.3.1. Necessity of $p \geq 2$. First we take the basic ‘‘vertical tube’’ solution. Let

$$F_0(x, y, z) = e^{-\frac{|x|^2 + |y|^2 + |z|^2}{2CR}}$$

so that

$$\begin{aligned} e^{\frac{it(\Delta_x + \Delta_y - \Delta_z)}{2}} F_0 &:= F(t, x, y, z) \\ &= \frac{1}{\left(1 + \frac{it}{CR}\right)^n \left(1 - \frac{it}{CR}\right)^{\frac{n}{2}}} e^{-\frac{|x|^2 + |y|^2}{2(CR+it)}} e^{-\frac{|z|^2}{2(CR-it)}}. \end{aligned} \quad (16)$$

The solution $F(t, x, y, z)$ is essentially 1 if $|(x, y, z)| \leq R^{1/2}$, $0 \leq t \leq R$. Also, the Fourier transform (in space) of F is essentially supported at frequencies $|(\xi, \eta, \zeta)| \leq R^{-1/2}$.

Pick roughly $R^{\frac{3n-1}{2}}$ points (x_k, y_k, z_k) which are spaced at distance $\sim R^{1/2}$ from each other on the surface $\{(x, y, z) : |x|^2 + |y|^2 = |z|^2, \frac{R}{2} \leq |x|, |y| \leq R\}$. Define $(\xi_k, \eta_k, \zeta_k) = \frac{1}{R}(x_k, y_k, z_k)$ so that

$$|\xi_k|^2 + |\eta_k|^2 = |\zeta_k|^2 \quad \text{and} \quad |(\xi_k, \eta_k, \zeta_k)| \sim 1.$$

Take the following initial conditions

$$G_0(x, y, z) = \sum e^{i(x \cdot \xi_k + y \cdot \eta_k - z \cdot \zeta_k)} F_0(x + x_k, y + y_k, z + z_k)$$

so that the solution is

$$\begin{aligned} G(t, x, y, z) \\ = \sum e^{i(x \cdot \xi_k + y \cdot \eta_k - z \cdot \zeta_k)} F(t, x + x_k - t\xi_k, y + y_k - t\eta_k, z + z_k - t\zeta_k), \end{aligned}$$

since $|\xi_k|^2 + |\eta_k|^2 = |\zeta_k|^2$.

Since the $\sim R^{\frac{3n-1}{2}}$ terms in G_0 are essentially orthogonal and each has L^2 norm $\sim R^{3n/4}$, we get

$$\|G_0\|_{L^2(dx dy dz)} \lesssim R^{\frac{3n}{2} - \frac{1}{4}}.$$

Moreover, since $|(\xi_k, \eta_k, \zeta_k)| \sim 1$, there also holds

$$\|G_0\|_{H^s(dx dy dz)} \lesssim R^{\frac{3n}{2} - \frac{1}{4}}. \quad (17)$$

From the expression of G , we see that

$$|G(t, x, y, z)| \gtrsim R^{\frac{3n-1}{2}} \quad \text{for } |(x, y, z)| \leq \frac{1}{100}, R - R^{\frac{1}{2}} < t < R.$$

Therefore,

$$\|G(t, x, x, x)\|_{L^p(dt)L^q(dx)} \gtrsim R^{\frac{1}{2p}} R^{\frac{3n-1}{2}}.$$

Recalling (17), if

$$\|G(t, x, x, x)\|_{L^p(dt)L^q(dx)} \lesssim \|G_0(x, y, z)\|_{H^s(dx dy dz)},$$

then $p \geq 2$. From a similar argument as in subsection 2.1.1, $p \geq 2$ is also necessary for estimates of the form

$$\|\nabla_x^\alpha G(t, x, x, x)\|_{L^p(dt)L^q(dx)} \lesssim \|G_0(x, y, z)\|_{H^s(dx dy dz)}.$$

2.3.2. Necessity of $q \geq 2$. Let $F(t, x, y, z)$ be the basic vertical tube solution of height R (as (16)). Let $m \gg 1$. Choose roughly $R^{mn - \frac{n}{2}}$ points x_k which are spaced at distance $\sim R^{\frac{1}{2}}$ in a large ball $B(0, R^m)$ of radius R^m in \mathbb{R}^n . Fix a unit vector $\xi \in S^{n-1}$.

We take initial conditions

$$G_0(x, y, z) = e^{i(x+y-z) \cdot \xi} \sum F_0(x + x_k, y + x_k, z + x_k),$$

so that the solution is

$$\begin{aligned} G(t, x, y) \\ = e^{\frac{-it}{2}} e^{i(x+y-z) \cdot \xi} \sum F(t, x + x_k - t\xi, y + x_k - t\xi, z + x_k - t\xi). \end{aligned}$$

There are roughly $R^{mn-\frac{n}{2}}$ terms in the sum. The summands are essentially orthogonal and each term has L^2 norm $\sim R^{3n/4}$, thus

$$\|G_0\|_{L^2(dx dy dz)} \sim R^{\frac{n}{2} + \frac{mn}{2}}.$$

On the other hand, each $F(t, x + x_k - t\xi, y + x_k - t\xi, z + x_k - t\xi)$ is essentially 1 on a tube T_k of radius $R^{1/2}$ and length R in $3n + 1$ dimensions, and rapidly decaying out of T_k . Note that at $t = 0$, T_k is centered at $(0, -x_k, -x_k, -x_k)$. Moreover, these tubes T_k are in the same direction $(1, \xi, \xi, \xi)$ and hence disjoint. Therefore, $|G(t, x, y, z)| \gtrsim 1$ on the union of the tubes T_k . In particular, $|G(t, x, x, x)| \gtrsim 1$ for $0 \leq t \leq R$ and $x \in B(t\xi, R^m)$. Thus

$$\|G(t, x, x, x)\|_{L^p([0,1])L^q(dx)} \gtrsim R^{\frac{mn}{q}}$$

(with a similar estimate for $|\nabla|^\alpha G(t, x, x, x)$), while $\|G_0\|_{H^s(dx dy)} \sim R^{\frac{n}{2} + \frac{mn}{2}}$ and $m \gg 1$, so $q \geq 2$ is necessary.

REFERENCES

- [1] T. Chen and N. Pavlović, *Derivation of the cubic NLS and Gross-Pitaevskii hierarchy from many body dynamics in $d = 3$ based on spacetime norms*, Ann. H. Poincaré, **15** (2014), 543–588.
- [2] T. Chen, Y. Hong and N. Pavlović, *Global Well-Posedness of the NLS System for Infinitely Many Fermions*, Archive for rational mechanics and analysis, April 2017, Volume **224**, Issue 1, pp 91–123.
- [3] X. Chen and J. Holmer, *On the Klainerman-Machedon Conjecture of the Quantum BBGKY Hierarchy with Self-interaction*, Journal of the European Mathematical Society, 2016 **18**, 1161–120.
- [4] X. Chen and J. Holmer, *Correlation structures, Many-body Scattering Processes and the Derivation of the Gross-Pitaevskii Hierarchy*. Int. Math. Res. Not. 2016 (**10**), 3051–3110.
- [5] L. Erdős, B. Schlein and H. T. Yau, *Derivation of the cubic non-linear Schrödinger equation from quantum dynamics of many-body systems*. Invent. Math. **167**, 515–614 (2007).
- [6] L. Erdős, B. Schlein and H. T. Yau, *Derivation of the Gross-Pitaevskii equation for the dynamics of Bose-Einstein condensate*. Annals Math. **172**, 291–370 (2010).
- [7] R. L. Frank, M. Lewin, E. H. Lieb, and R. Seiringer, *Strichartz inequality for orthonormal functions*, J. Eur. Math. Soc., (2014).
- [8] R. L. Frank and J. Sabin, *Restriction theorems for orthonormal functions, Strichartz inequalities and uniform Sobolev estimates*, American Journal of Mathematics Johns Hopkins University Press Volume **139**, Number 6, December 2017 pp. 1649–1691.
- [9] M. Grillakis and M. Machedon, *Pair excitations and the mean field approximation of interacting Bosons, II*, Communications in PDE, Vol **42**, No 1, 24–67 (2017).

- [10] M. Grillakis and M. Machedon, *Uniform in N estimates for a Bosonic system of Hartree-Fock-Bogoliubov type*, Communications in PDE, Volume **44**, Number 12, 2019, pp. 1431–1465.
- [11] S. Klainerman and M. Machedon, *On the uniqueness of solutions to the Gross-Pitaevskii hierarchy*. Comm. Math. Phys. **279**, 169–185 (2008).

UNIVERSITY OF MARYLAND, COLLEGE PARK
E-mail address: `xdu@math.umd.edu`

UNIVERSITY OF MARYLAND, COLLEGE PARK
E-mail address: `mxm@math.umd.edu`